

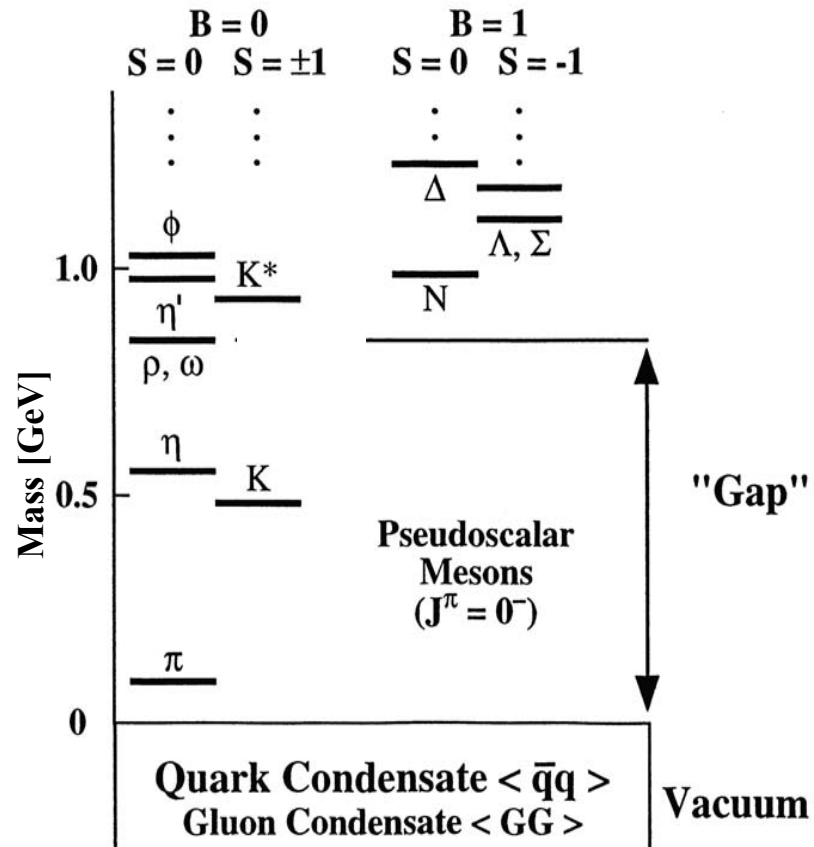
In-medium modifications of vector mesons in elementary reactions and heavy-ion collisions

**Volker Metag
II. Physikalisches Institut
Universität Giessen
Germany**

- **theoretical motivation and predictions**
- **medium modifications observed in elementary reactions:
KEK-E325, Jlab-G7, CBELSA/TAPS**
- **medium modifications observed in heavy-ion collisions:
DLS, HADES, CERES, NA60**
- **summary and conclusions**

**Quark Matter 2006
Shanghai, China, Nov 12th-20th 2006**

Motivation



- hadrons = excitations of the QCD vacuum
- QCD-vacuum: complicated structure characterized by condensates
- in the nuclear medium: condensates are changed
→ change of the hadronic excitation energy spectrum

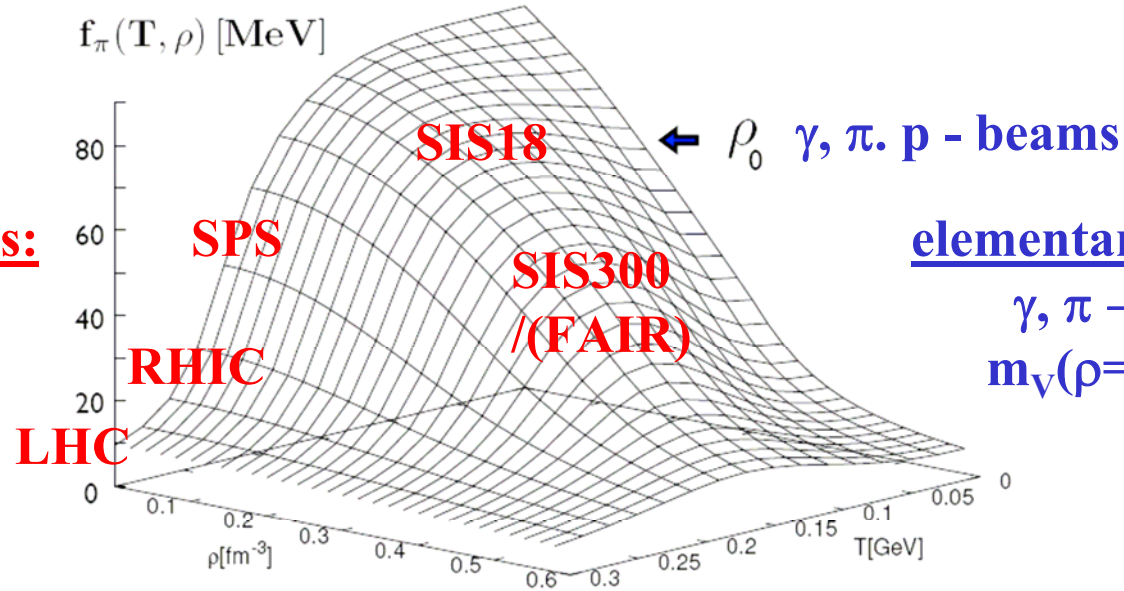
G.E.Brown and M. Rho, $\frac{m^*}{m} \approx \frac{\langle \bar{q}q \rangle^*}{\langle \bar{q}q \rangle} \approx 0.8$ ($\rho \approx \rho_0$)
PRL 66 (1991) 2720

T.Hatsuda and S. Lee, $\frac{m_V^*}{m_V} = \left(1 - \alpha \frac{\rho_B}{\rho_0} \right)$; $\alpha \approx 0.18$
PRC 46 (1992) R34

⇒ widespread theoretical and experimental activities
to study in-medium modifications of hadrons

chiral condensate as function of baryon density ρ_B and temperature T

C. Ratti, M. Thaler, W. Weise, PRD73 (2006) 014019



heavy ion reactions:



$$m_V(\rho \gg \rho_0; T \gg 0)$$

elementary reaction:



$$m_V(\rho = \rho_0; T = 0)$$

$$\frac{f_\pi^2(T, \rho)}{f_\pi^2(0)} \simeq \frac{\langle \bar{q}q \rangle_{T, \rho}}{\langle \bar{q}q \rangle_0} \simeq 1 - \frac{T^2}{8f_\pi^2} - \frac{\sigma_N}{m_\pi^2 f_\pi^2} \rho + \dots$$

$\langle q\bar{q} \rangle$ is not an observable!!

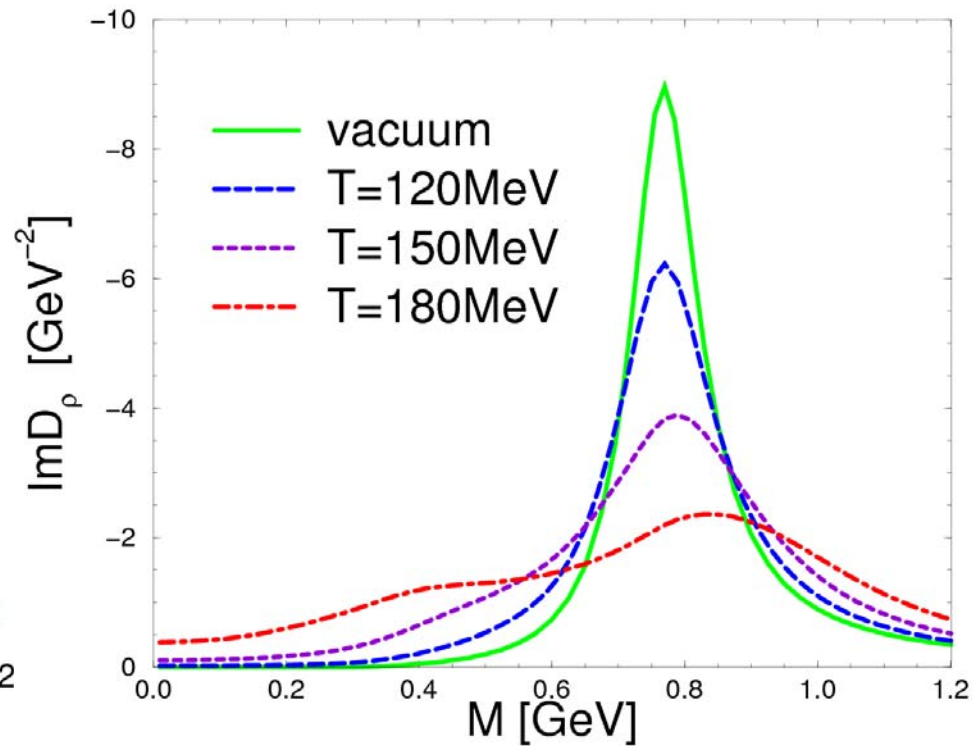
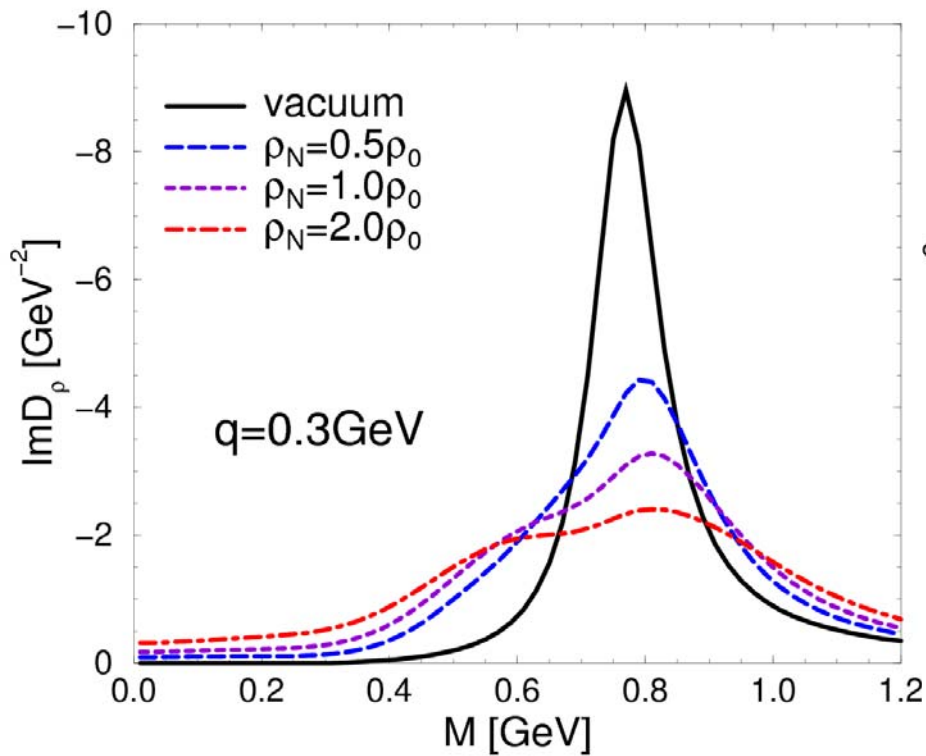
QCD sum rules: provide link between hadronic observables and condensates

$$\frac{Q^2}{24\pi^2} \int ds \frac{R(s)}{(s+Q^2)^2} = \frac{1}{16\pi^2} \left(1 + \frac{\alpha_s}{\pi} \right) + \frac{1}{Q^4} \left[m_q \langle \bar{q}q \rangle + \frac{1}{24} \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle \right] + \text{higher order terms}$$

hadronic spectral function: $R(s) \sim F^2 \frac{1}{\pi} \frac{\sqrt{s} \Gamma(s)}{(s - M_\rho^2)^2 + s(\Gamma(s))^2}$

ρ spectral function for different densities and temperatures

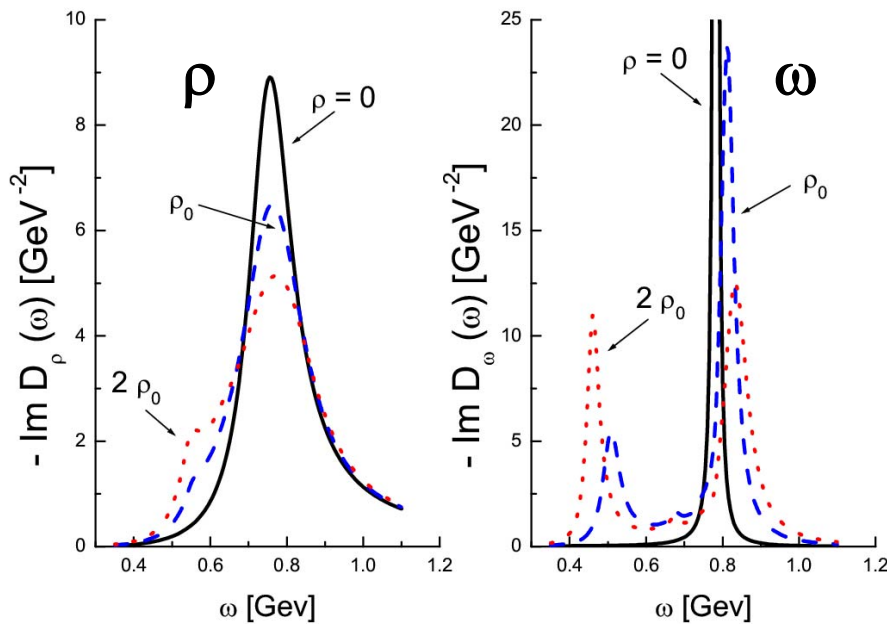
R. Rapp and J. Wambach, EPJA 6 (1999) 415



broadening and strength below vacuum pole mass

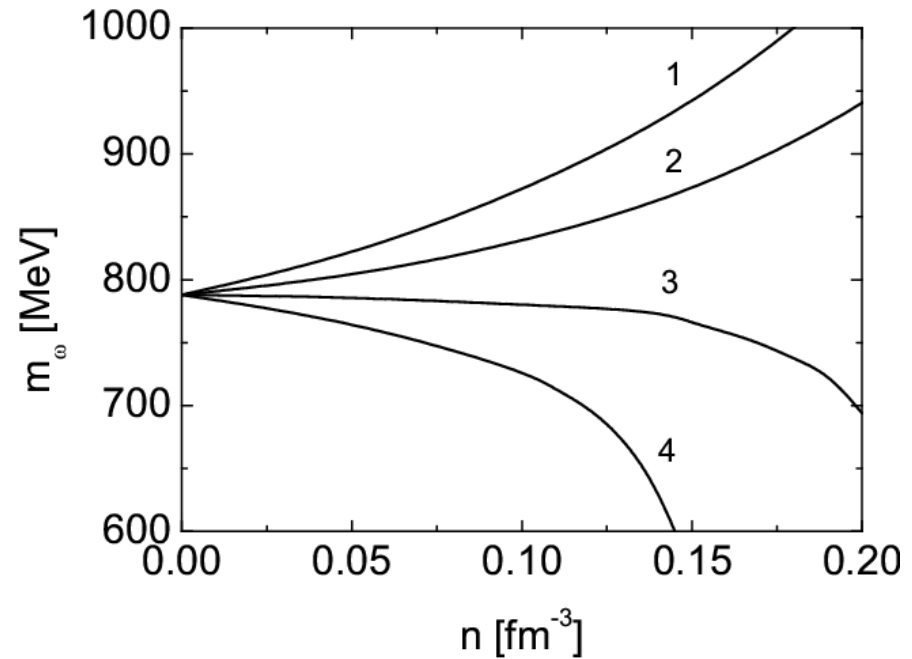
Model predictions for spectral functions of ρ and ω mesons

M. Lutz et al., Nucl. Phys. A 706 (2002) 431



structure in spectral function due to coupling to baryon resonances

S. Zchocke et al., Phys. Lett. B 562 (2003) 562



variation in ω -mass due to density dependence of 4-quark condensate

experimental approach: dilepton spectroscopy: $\rho, \omega, \phi \rightarrow e^+e^-$

reconstruction of invariant mass from 4-momenta of decay products:

$$m_\omega(\rho, T, \vec{p}) = \sqrt{(\mathbf{p}_1 + \mathbf{p}_2)^2}$$

essential advantage: no final state interactions !!

I. Information on medium modifications of mesons from elementary reactions

advantage: well controlled conditions:

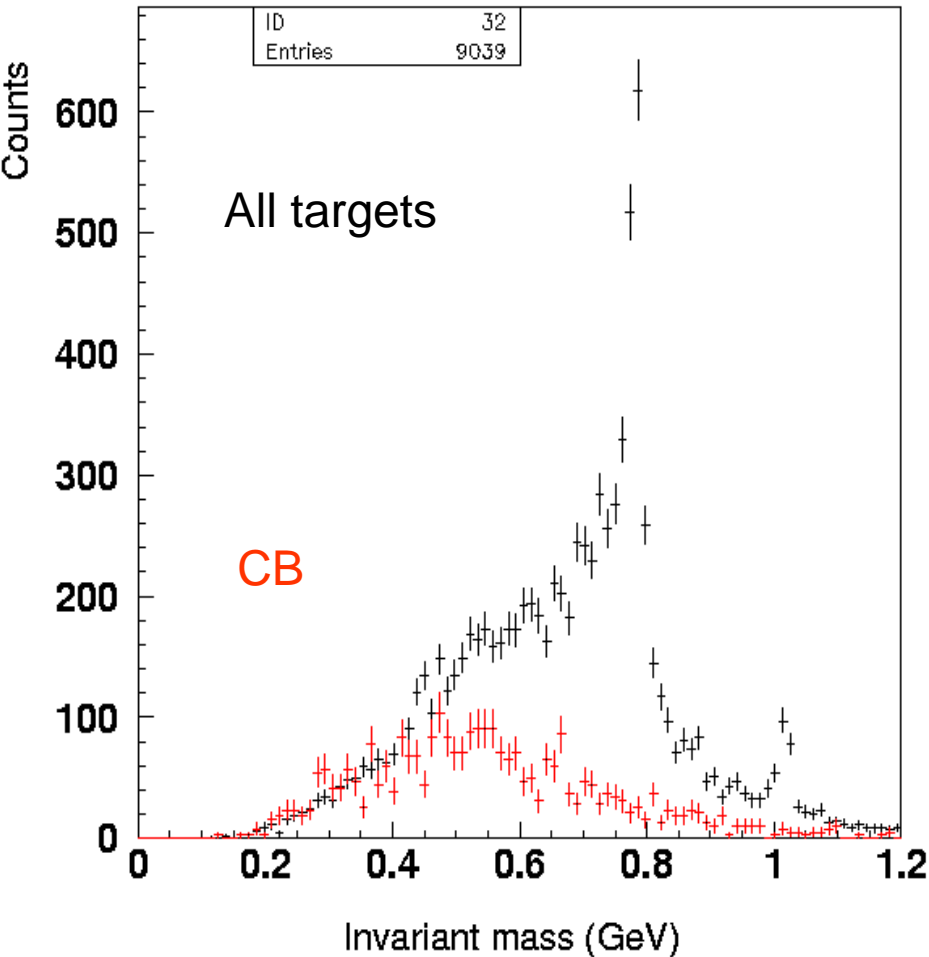
important for theoretical interpretation

no time dependence of baryon density $\rho_B \neq \rho_B(t); T=0;$

disadvantage: small medium effects since $\rho \leq \rho_0$ and $T=0$

JLAB-CLAS: G7 $\gamma A \rightarrow e^+e^- + X$

$E_\gamma = 0.6-3.8$ GeV

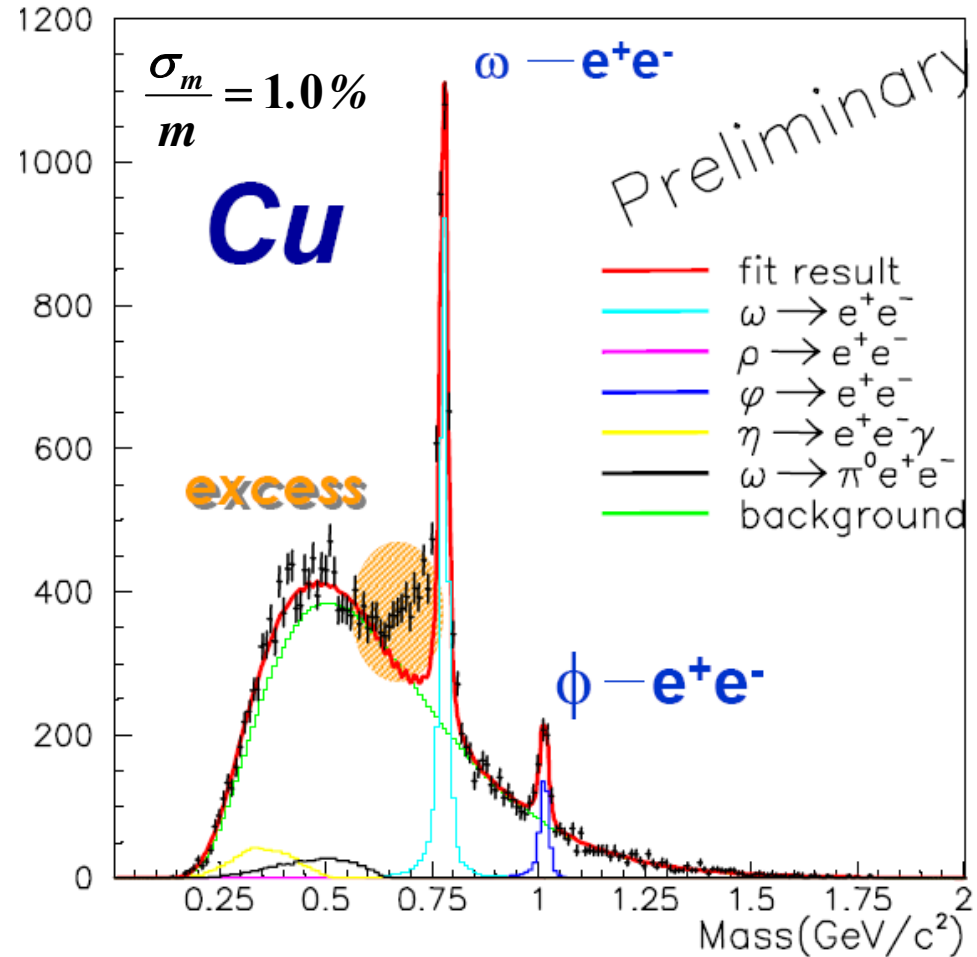


s. plenary talk by Ch. Djalali

ω, ϕ -peaks observed;

additional yield above combinatorial background assigned to $\rho \rightarrow e^+e^-$

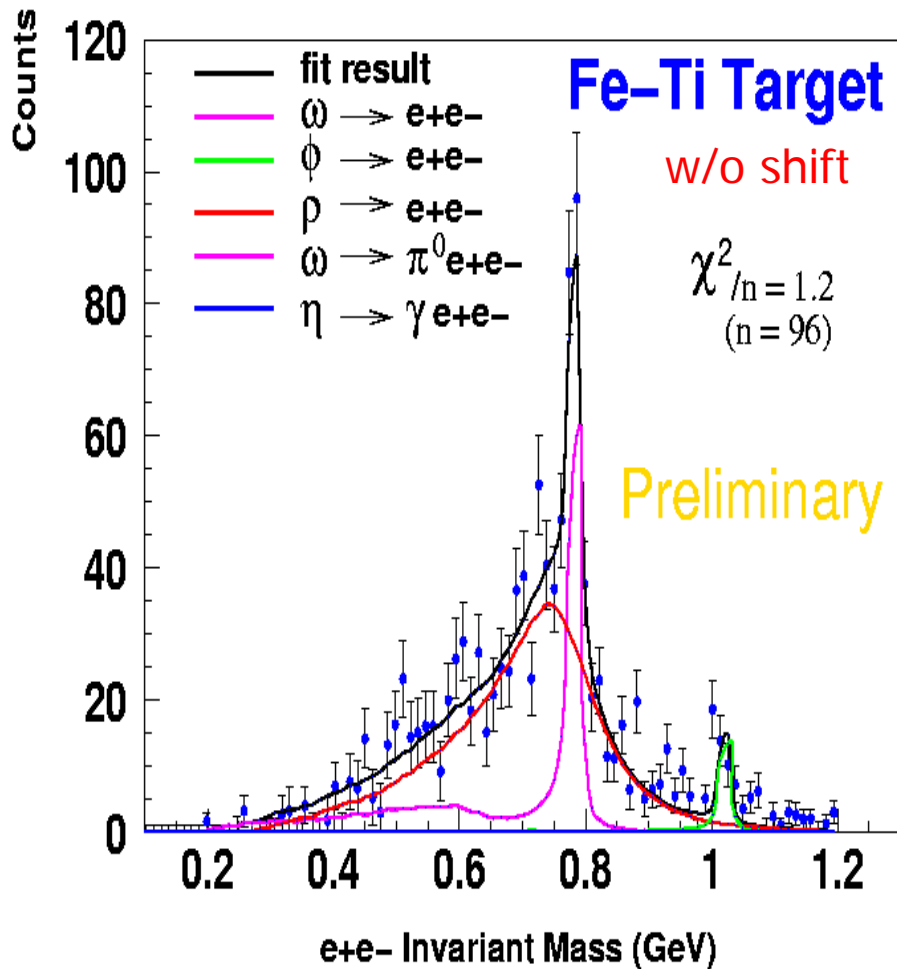
KEK-E325: p (12 GeV) $A \rightarrow \rho, \omega + X$



s. talk by M. Naruki

comparison of JLAB and KEK results

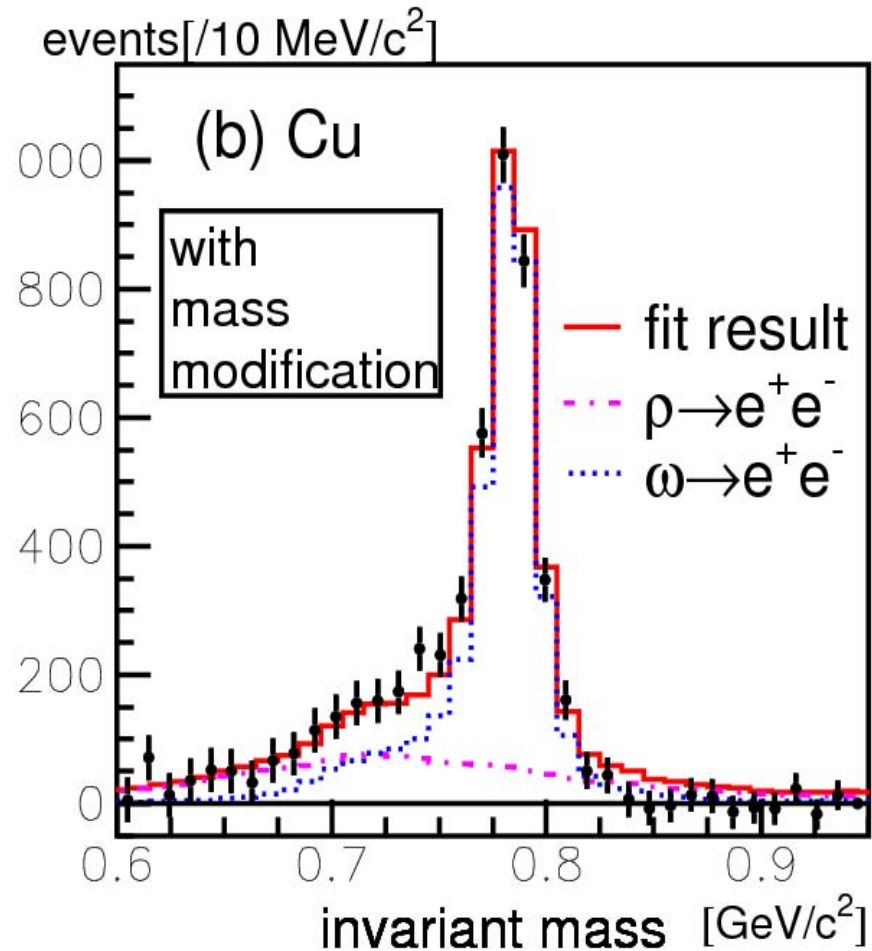
JLAB-CLAS: G7 $\gamma A \rightarrow e^+e^- + X$



ρ slightly broadened; no mass shift

No consistent picture!!

KEK-E325: p (12 GeV) A $\rightarrow \rho, \omega + X$



ρ shifted in mass: $m_\rho = m_0 \left(1 - 0.092 \frac{\rho}{\rho_0} \right)$;
no broadening!!

M. Naruki et al., PRL 96 (2006) 092301

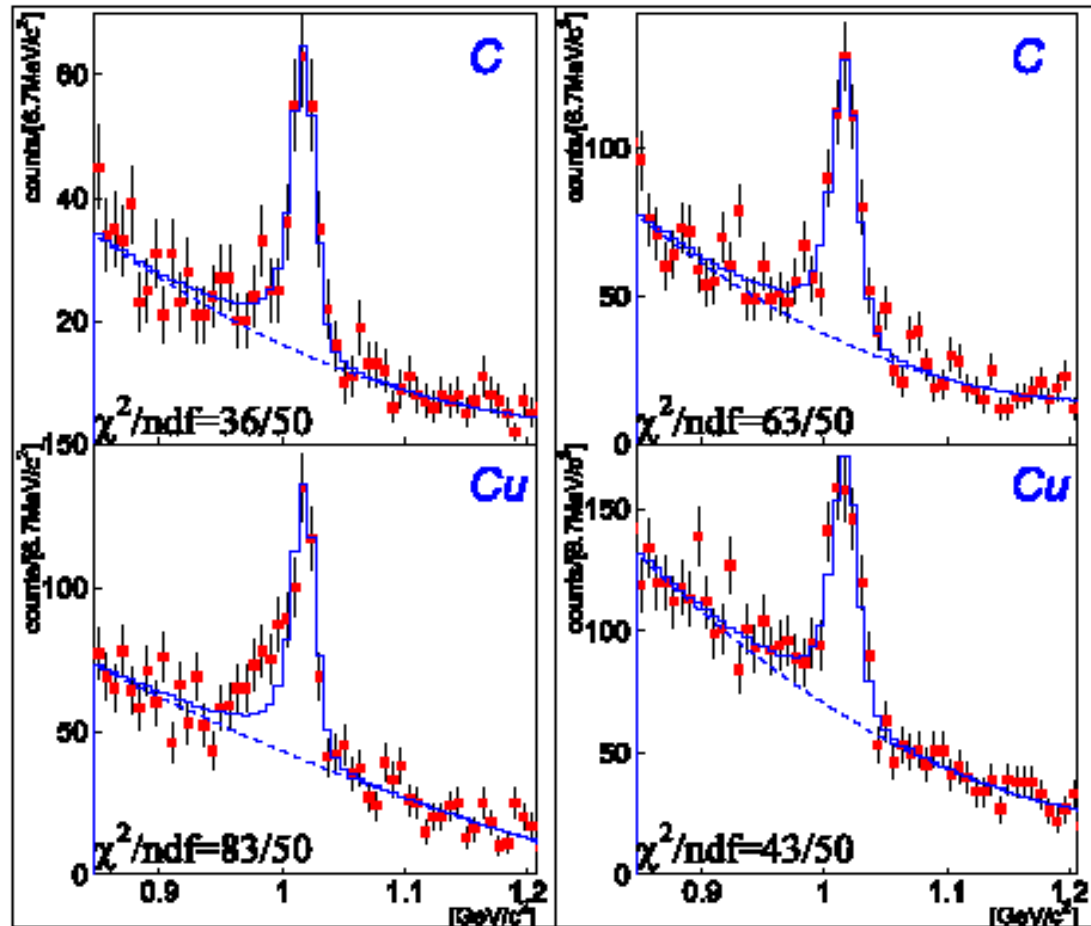
Φ - meson in the nuclear medium

R.Muto et al.,
nucl-ex/0511019

KEK-E325: p (12 GeV) $\Lambda \rightarrow \rho, \omega + X; \Phi \rightarrow e^+e^-$

$\beta\gamma < 1.25$ (Slow)

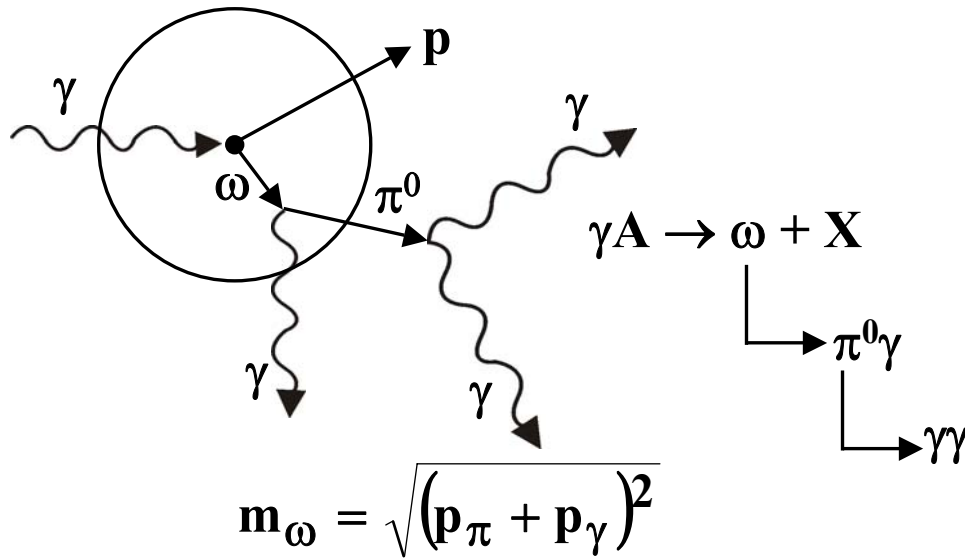
$1.25 < \beta\gamma < 1.75$



mass shift of Φ - meson for low recoil momenta in Cu: $m_\Phi = m_0 \left(1 - 0.04 \frac{\rho}{\rho_0} \right)$;

ω -mass in nuclei from photonuclear reactions

J.G.Messchendorp et al., Eur. Phys. J. A 11 (2001) 95



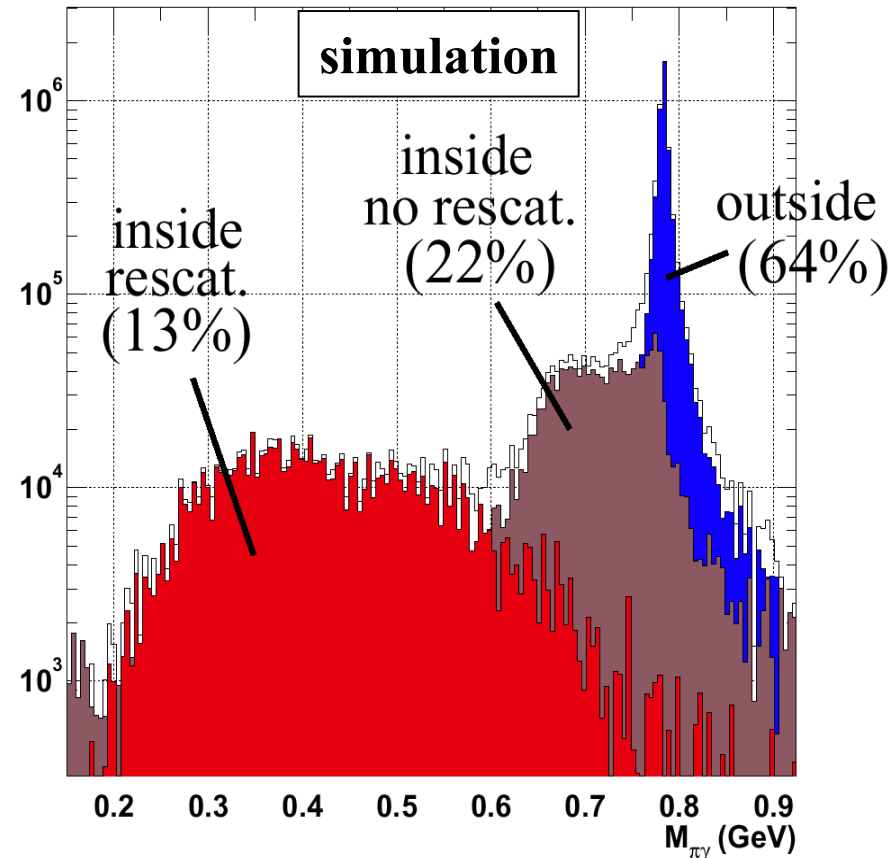
advantage:

- $\pi^0 \gamma$ large branching ratio (8 %)
- no ρ -contribution ($\rho \rightarrow \pi^0 \gamma : 7 \cdot 10^{-4}$)

disadvantage:

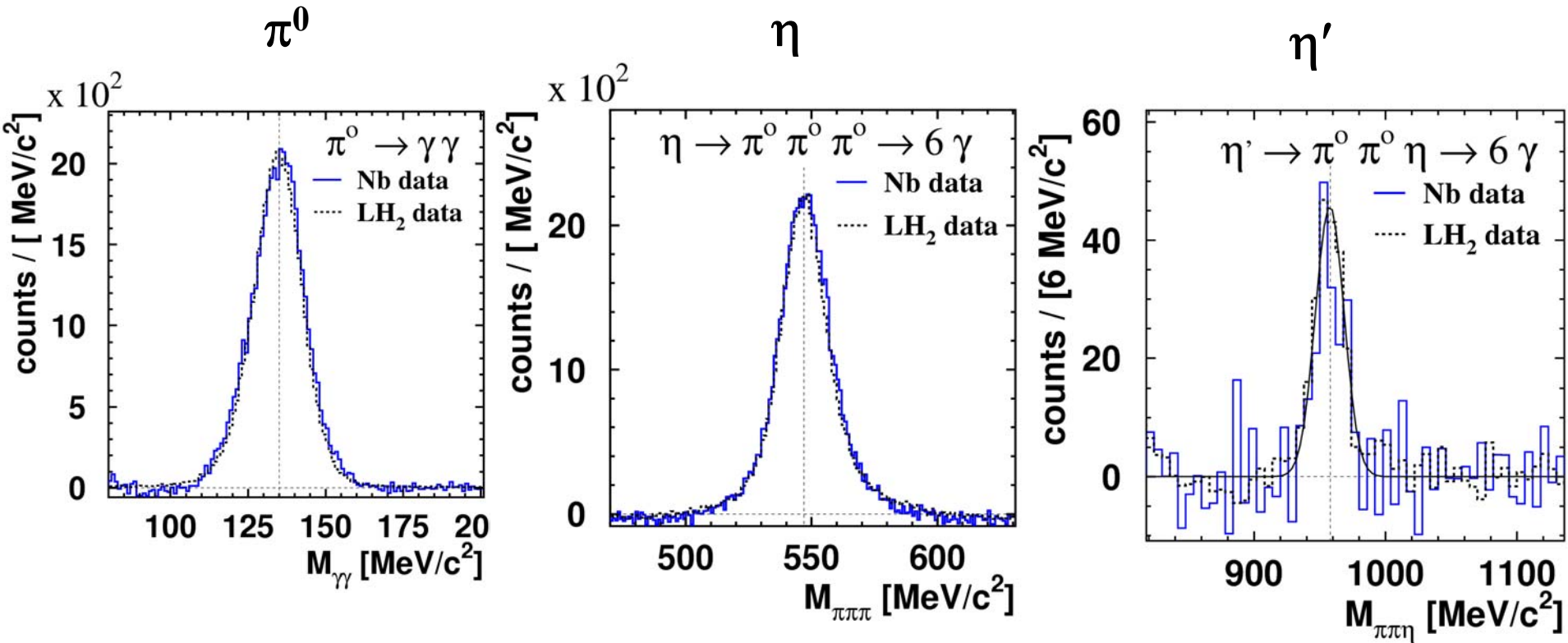
- π^0 -rescattering

$\gamma + \text{Nb} @ 1.2 \text{ GeV}$



no distortion by pion rescattering
 expected in mass range of interest;
 further reduced by requiring $T_{\pi} > 150 \text{ MeV}$

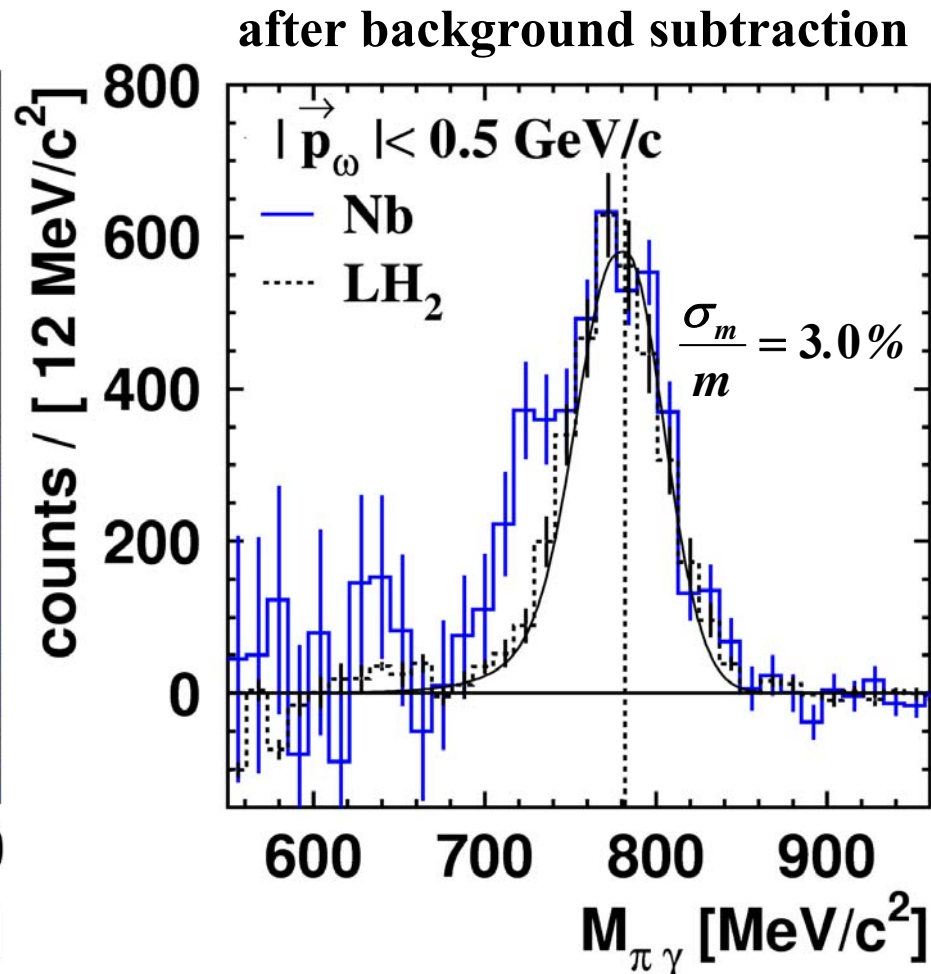
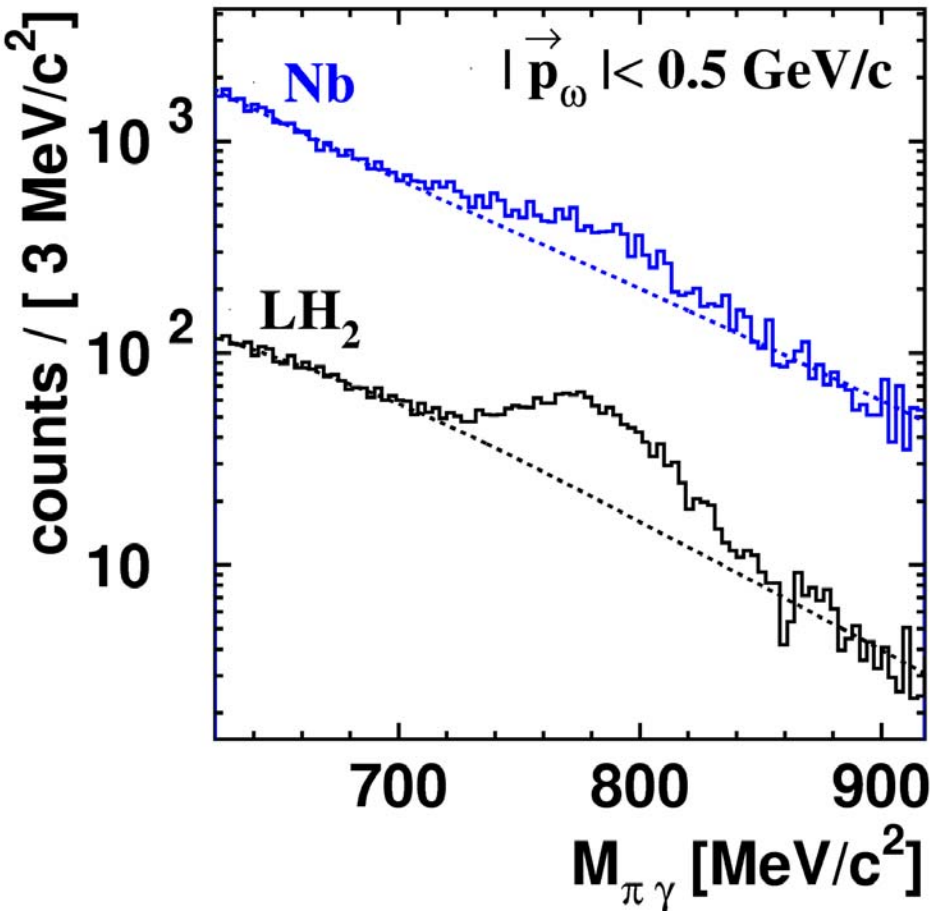
comparison of meson masses and lineshapes for LH₂ and nuclear targets after background subtraction



No change of mass and lineshape for longlived mesons (π^0 , η , η') decaying outside nuclei

inclusive $\omega \rightarrow \pi^0 \gamma$ signal for LH₂ and Nb target

D. Trnka et al., PRL 94 (2005) 192203



difference in line shape of ω signal for proton and nuclear target
consistent with $m_\omega = m_0 (1 - \alpha \rho/\rho_0)$ for $\alpha = 0.13$

refined analysis requiring recoil proton and p- ω coplanarity

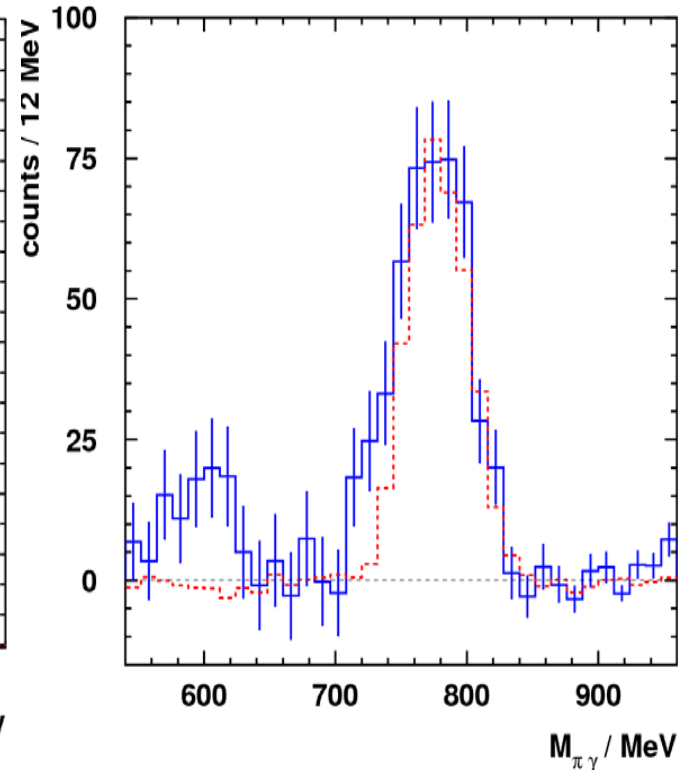
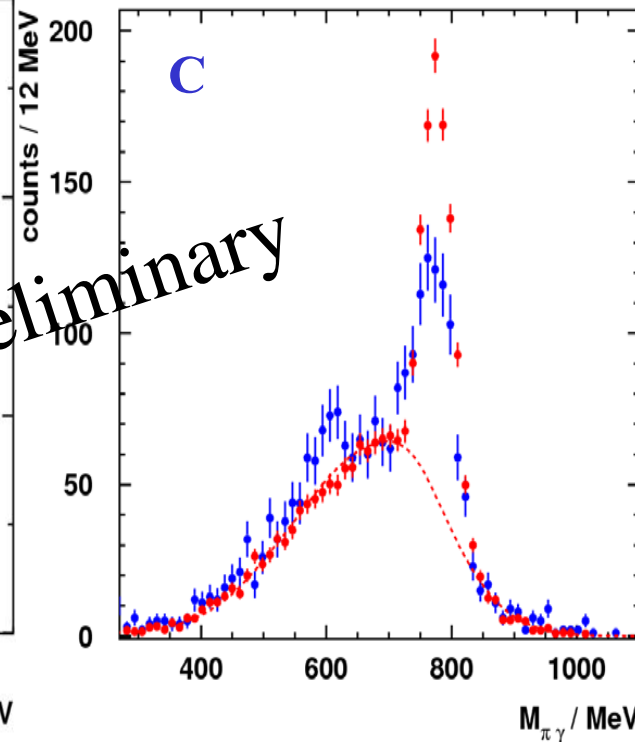
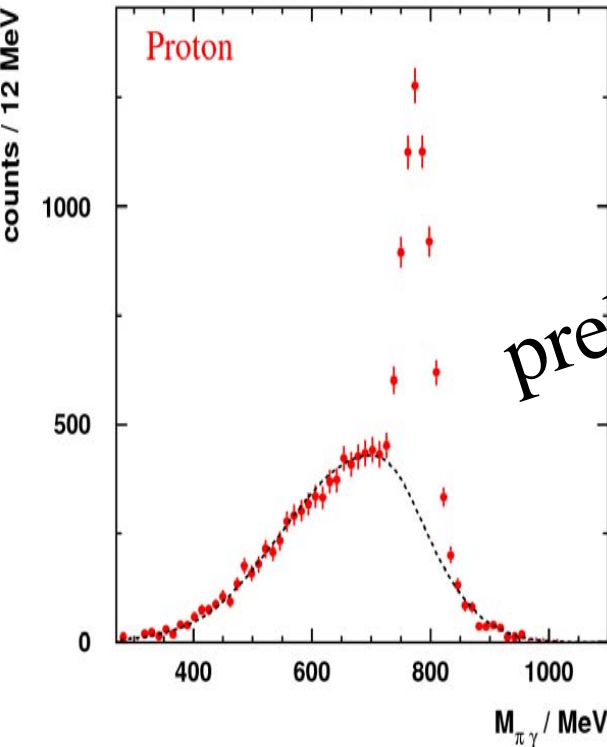
D. Trnka (Gießen) priv. com.

No background subtraction!!!

$|p_\omega| < 500$ MeV/c

$|p_\omega| < 500$ MeV/c

after LH₂ background subtraction



\Rightarrow difference in ω - line shape for proton and nuclear target confirmed;
no upward mass shift of ω meson!

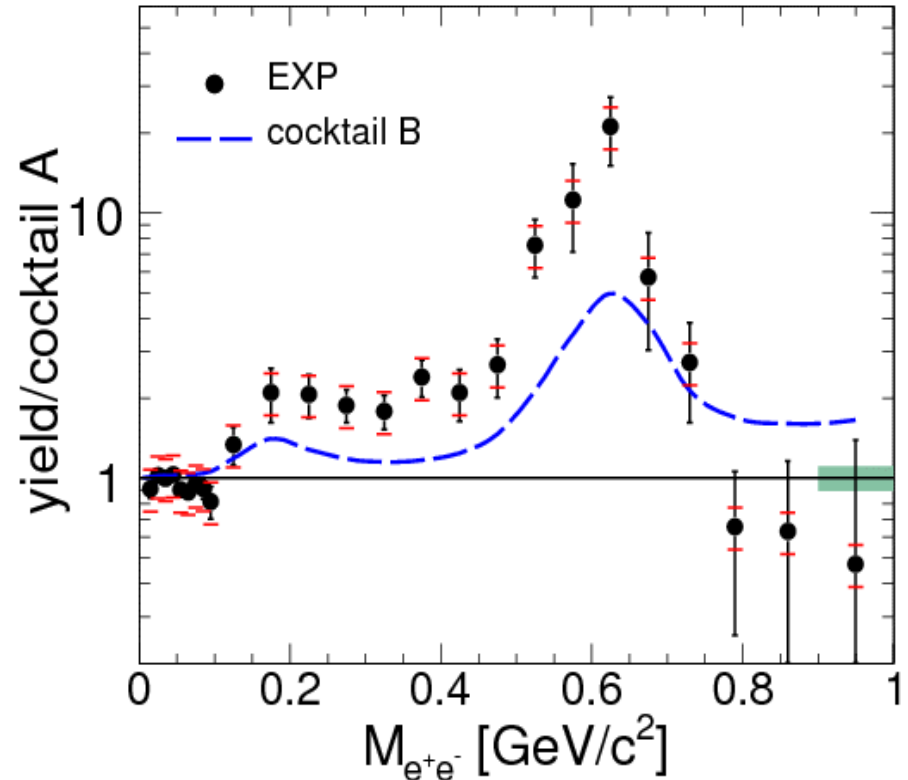
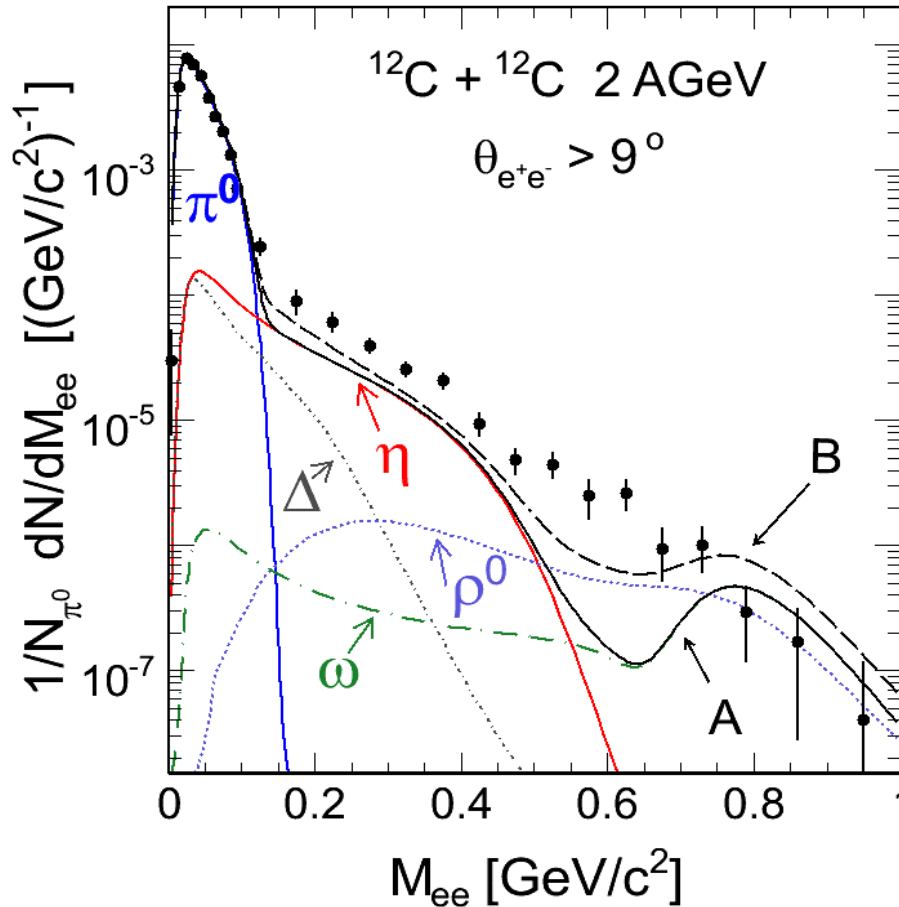
\Rightarrow additional structure at ≈ 600 MeV!! (also seen for heavier targets)
fragmentation of ω strength or background ??? under investigation

II. Information on medium modifications of mesons from heavy-ion collisions

advantage: sizable effects due to high densities and temperatures

**disadvantage: any signal represents an integration over the full
space-time history of the heavy-ion collision with
strong variations in densities and temperatures**

G. Agakichiev et al, subm. to PRL



sizable excess in dilepton yield
 relative to contributions from
 long-living mesons

origin of excess-yield??

cocktail A: post freeze-out hadronic sources:

$\pi^0, \eta \rightarrow \gamma e^+e^-$; $\omega \rightarrow \pi^0 e^+e^-$; $\omega \rightarrow e^+e^-$;

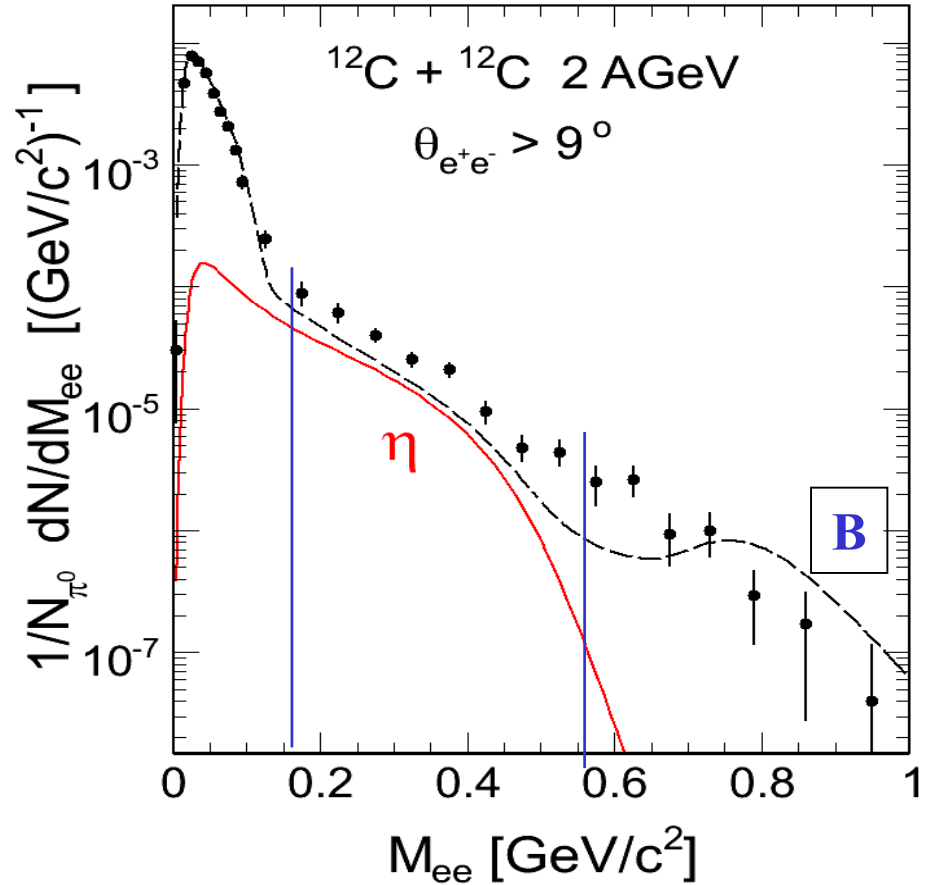
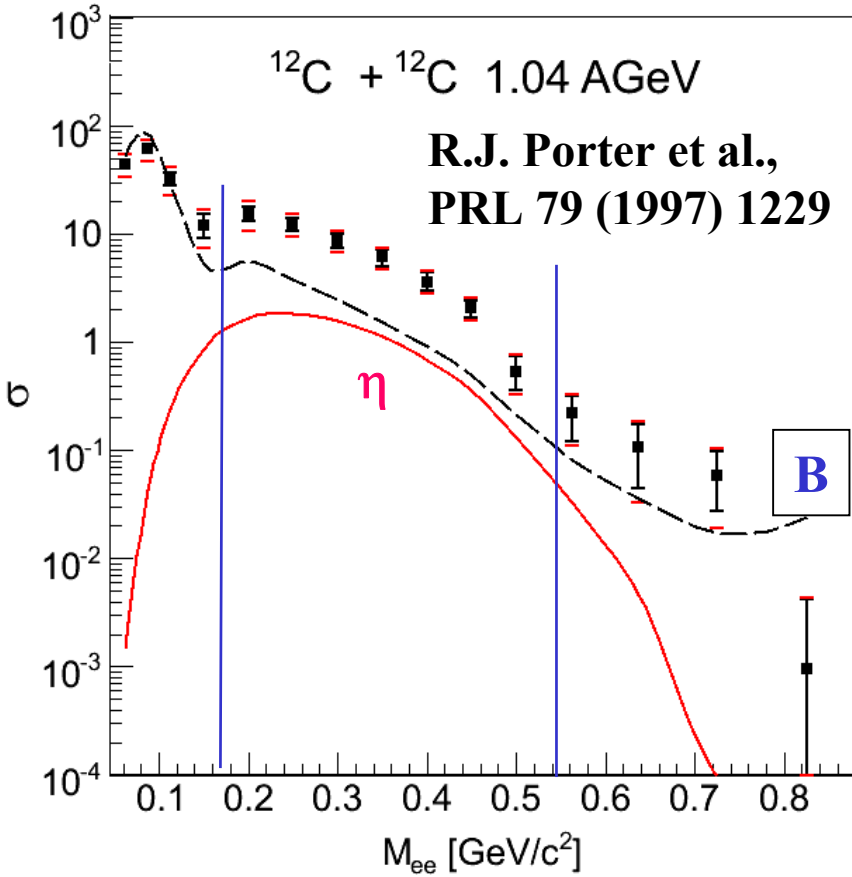
cocktail B: additional fireball sources:

$\Delta \rightarrow N e^+e^-$; $\rho \rightarrow e^+e^-$;

Dilepton production in C+C

DLS

HADES



for $170 \text{ MeV} < M_{ee} < 550 \text{ MeV}$

$$F(1.04 \text{ AGeV}) = 6.5 \pm 0.5(\text{stat}) \pm 2.1(\text{sys})$$

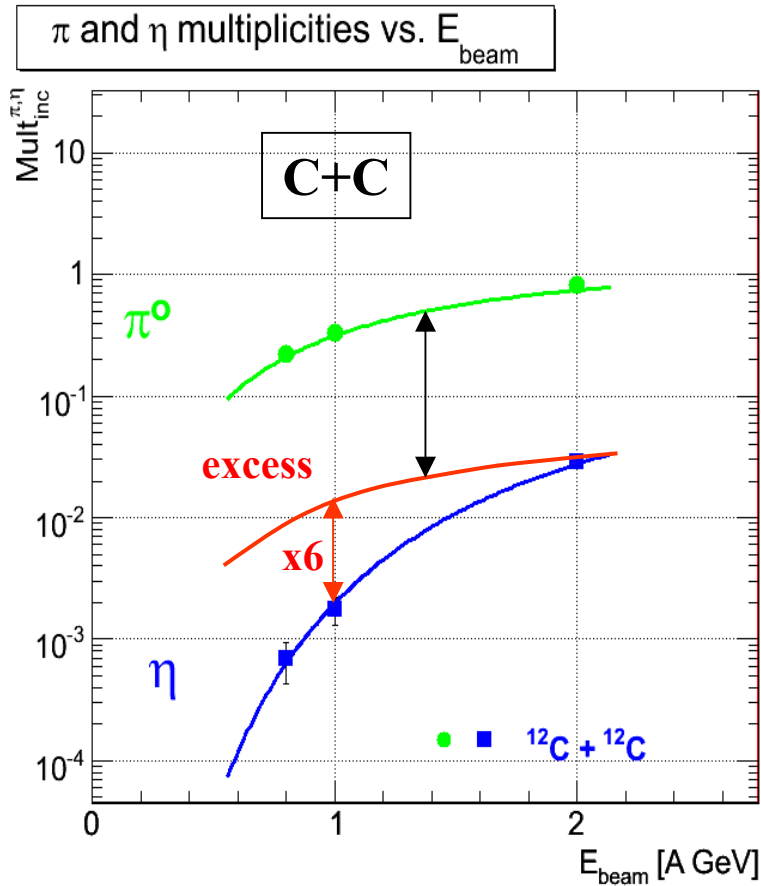
$$F(2.0 \text{ AGeV}) = 2.07 \pm 0.21(\text{stat}) \pm 0.38(\text{sys})$$

$$F = \frac{Y_{tot}}{Y_\eta} = \frac{Y_{exc} + Y_\eta}{Y_\eta} = \frac{Y_{exc}}{Y_\eta} + 1$$

$$\Rightarrow Y_{exc} = (F - 1)Y_\eta$$

comparison of excess yield with π^0 , η - cross section

TAPS data: { R. Averbeck et al., Z. Phys. A359 (1997) 65
R. Holzmann et al., PRC 56 (1997) R2920



$$\frac{Y_{exc}(2.0 \text{ A GeV})}{Y_{exc}(1.04 \text{ A GeV})} = \frac{(F(2.0) - 1)}{(F(1.04) - 1)} \cdot \frac{Y_{\eta}(2.0)}{Y_{\eta}(1.04)}$$

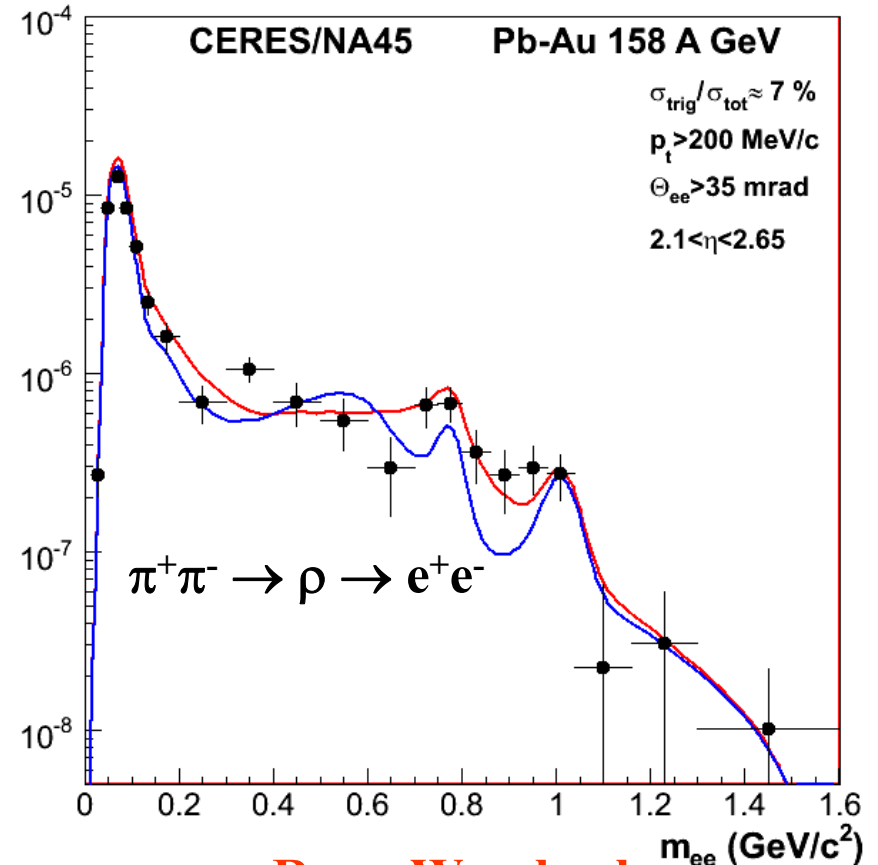
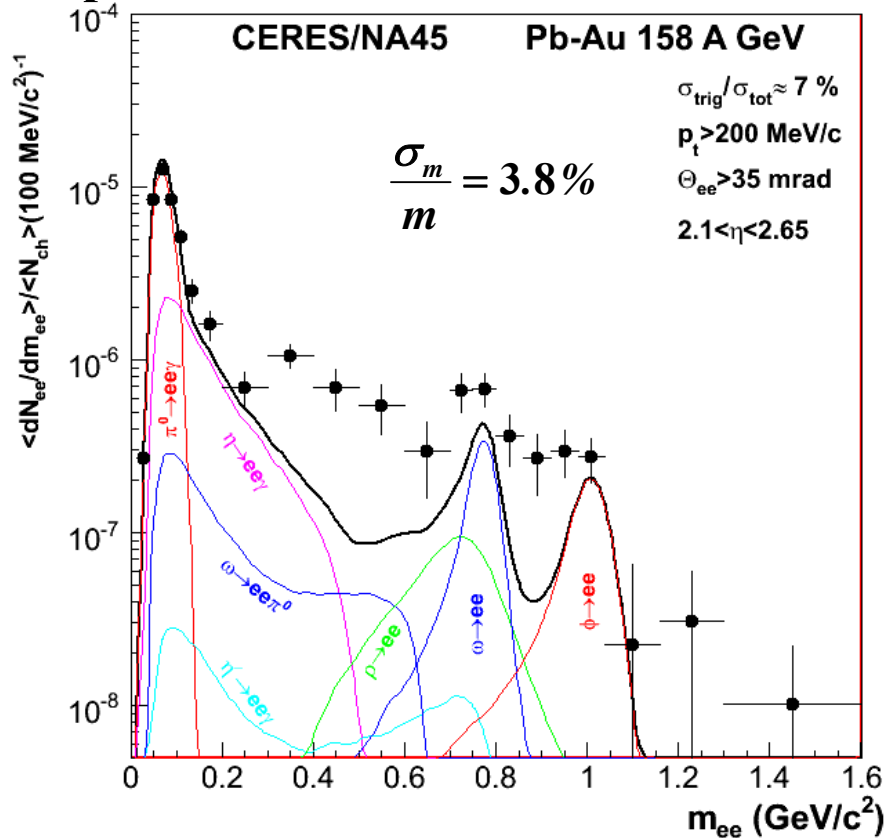
$$= 2.5 \pm 0.5(\text{stat}) \pm 1.5(\text{sys})$$

$$\frac{Y_{\pi}(2.0 \text{ A GeV})}{Y_{\pi}(1.04 \text{ A GeV})} = 2.3 \pm 0.3$$

excess yield scales with π^0 cross section !! and amounts to $\approx 2 \cdot 10^{-5}$ of the π^0 -yield (assuming similar acceptances for the excess- and η -yields)

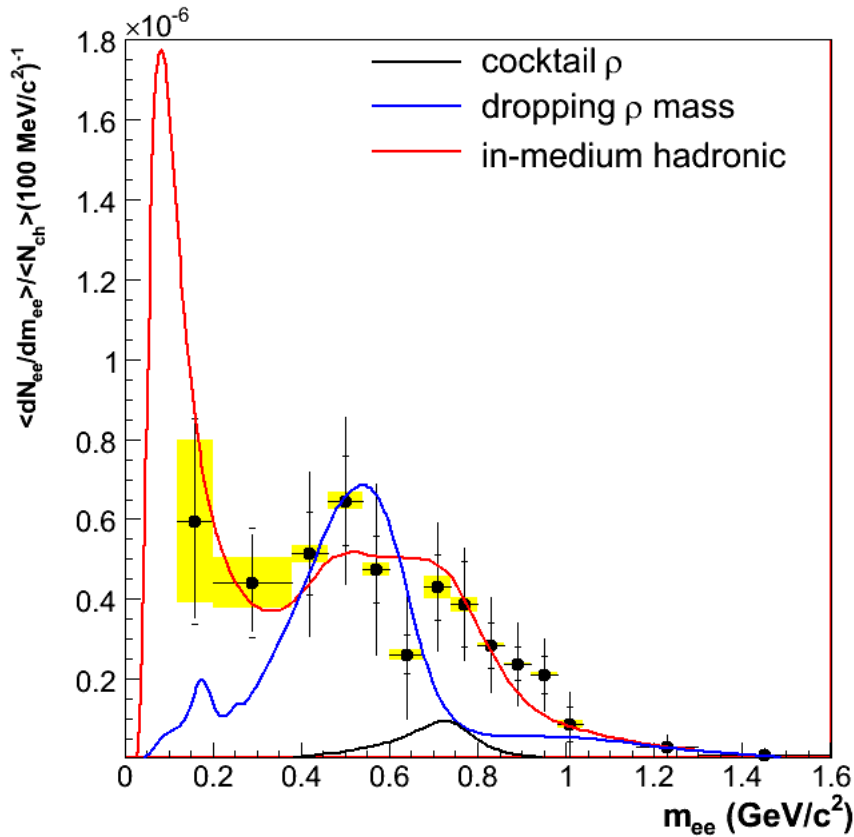
\Rightarrow dilepton excess associated with excitation of baryon resonances !!

data in comparison to
post freeze out hadronic cocktail

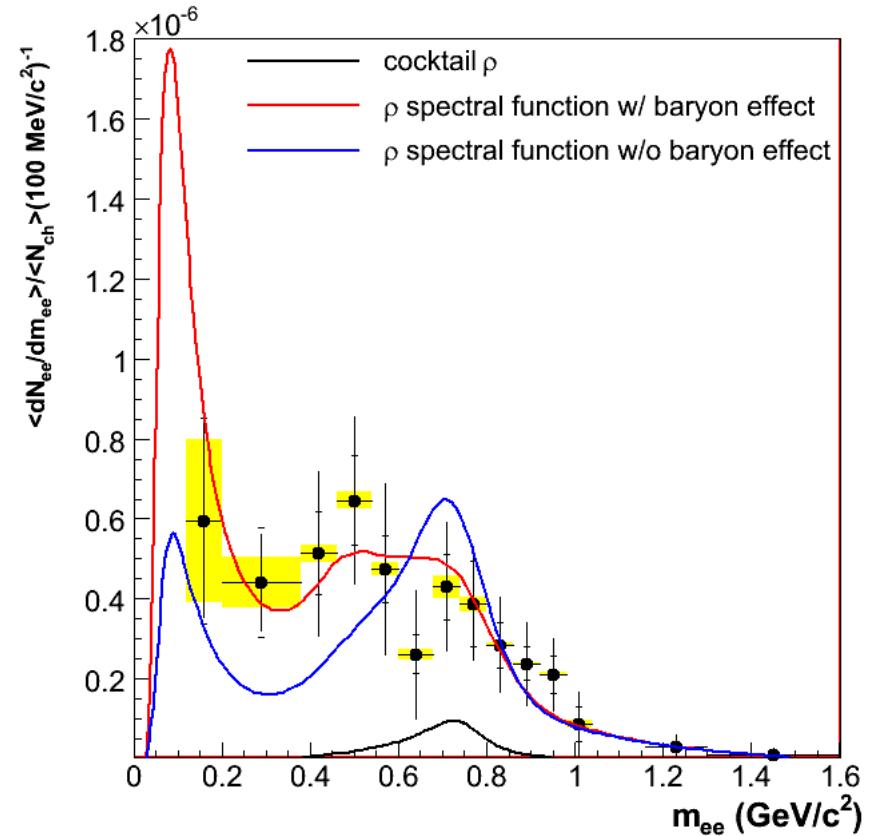


for $0.2 < m_{ee} < 1.1 \text{ GeV}/c^2$
excess over hadronic decay contribution:
 $2.45 \pm 0.30(\text{stat}) \pm 0.38(\text{syst}) \pm 0.74(\text{decays})$

— Rapp Wambach
— broadened ρ spect. function
— Brown Rho
— dropping ρ mass

CERES e^+e^- pair yield after subtraction of hadronic cocktail excluding ρ 

excess smeared out over wider mass range than expected for dropping ρ mass scenario



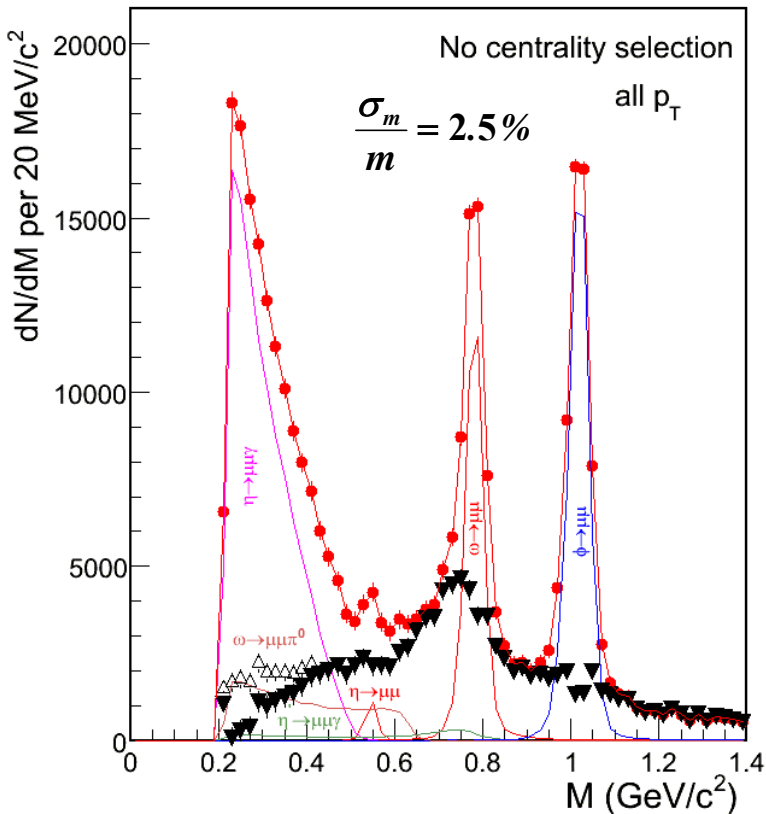
strength of di-lepton yield at low masses due to coupling to baryons

(similar to the observation by DLS, HADES)

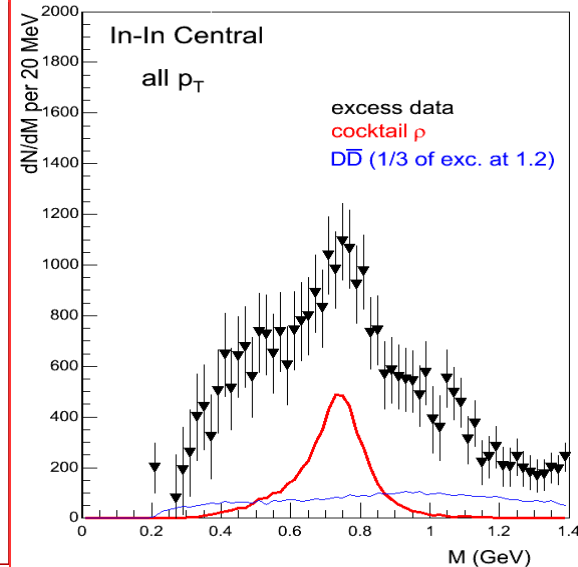
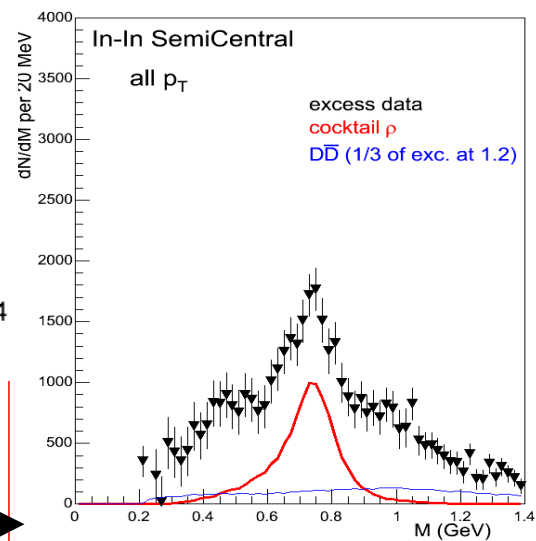
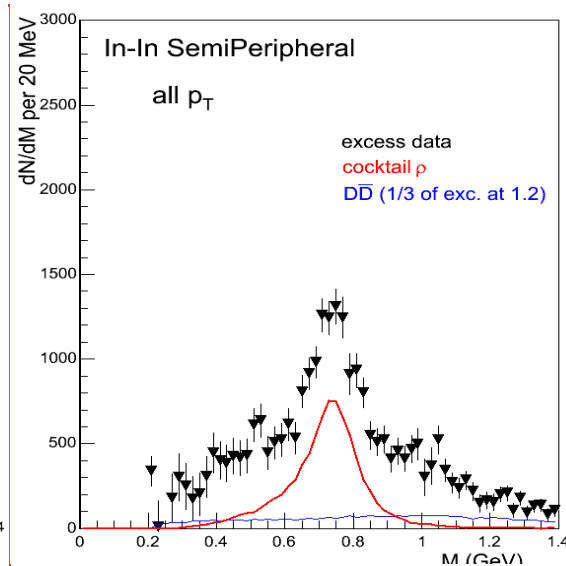
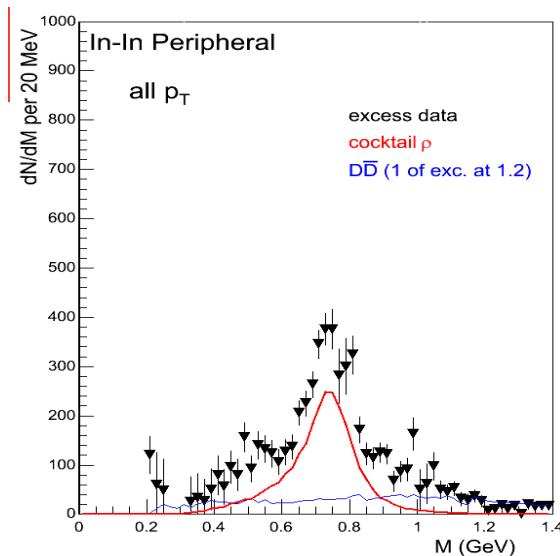
R. Arnaldi et al., PRL 96 (2006) 162302

after subtraction of hadronic cocktail

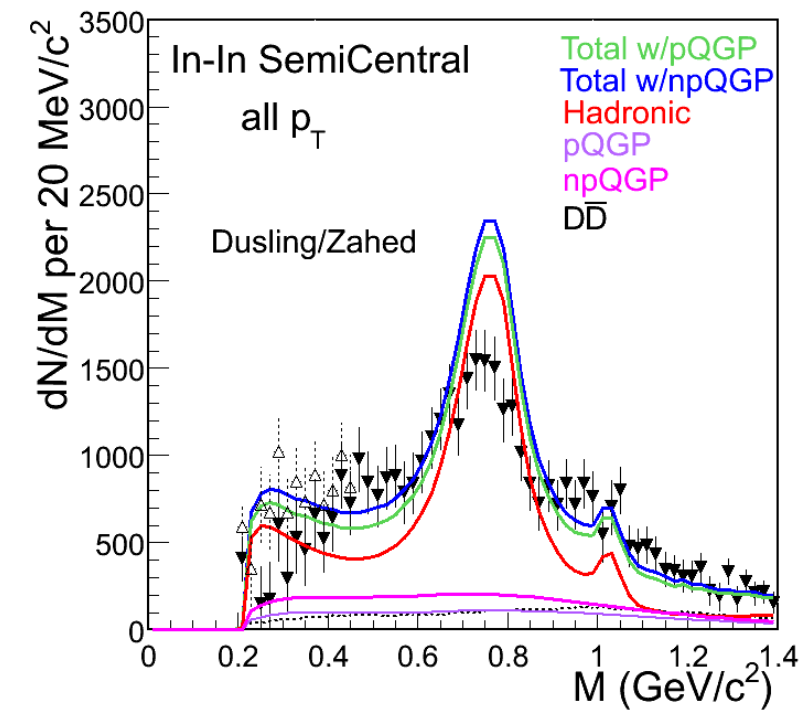
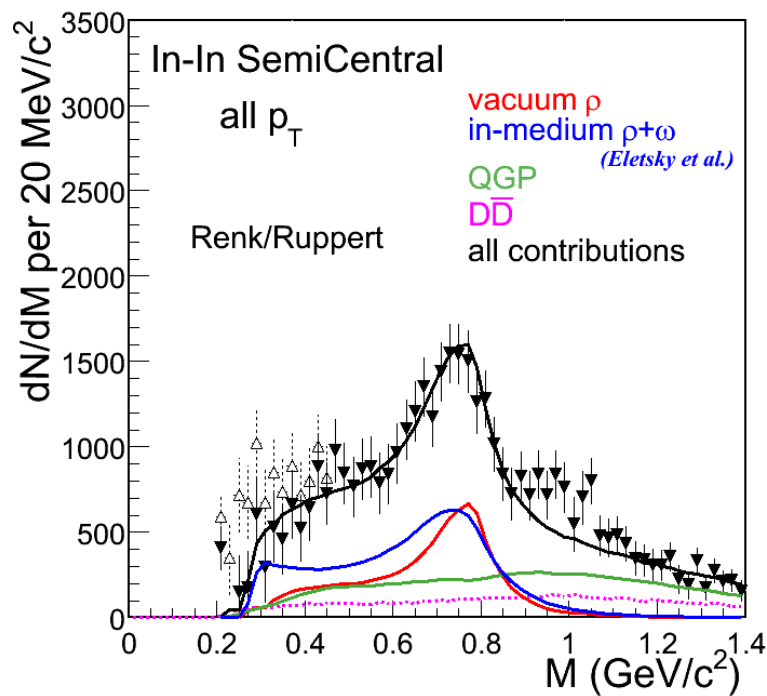
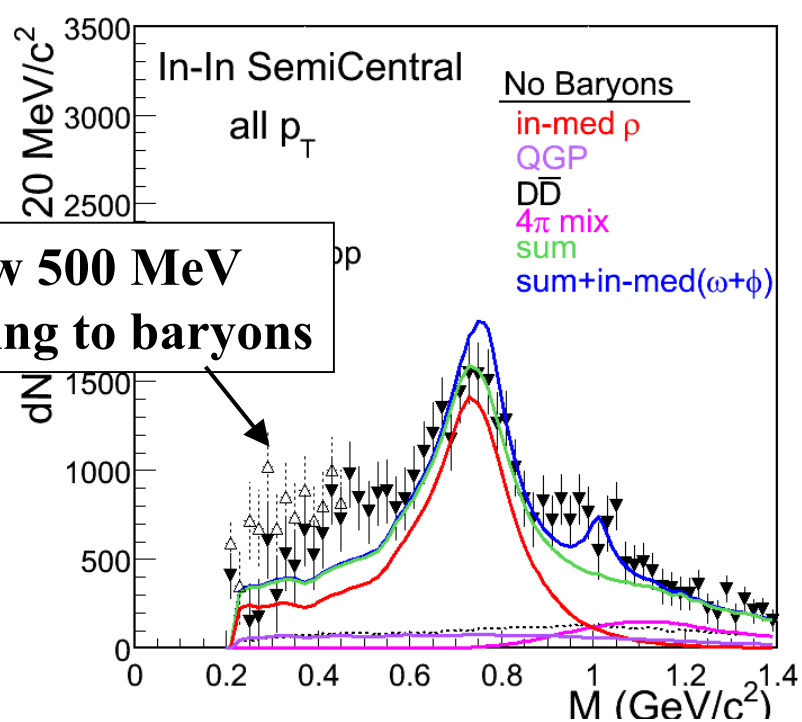
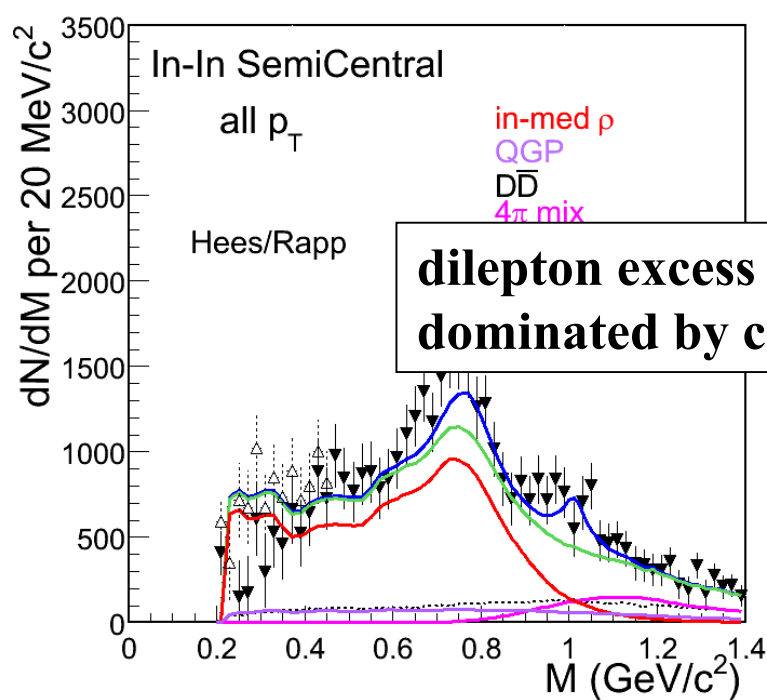
isolation of $\mu^+ \mu^-$ excess



excess increases with centrality



\Rightarrow broadening of ρ meson; no mass shift



summary

current status of in-medium modifications of vector mesons

	KEK	Jlab	CBELSA/TAPS	CERES	NA 60
ω	–	–	mass shift: -14% $\Gamma_{\omega}(\rho=\rho_0)\approx 100\text{MeV}$	–	–
ρ	mass shift: -9% no broadening	no mass shift some broadening	–	broadening favored over density dependent mass shift	no mass shift strong broadening
Φ	mass shift: -4% $\Gamma_{\phi}(\rho_0)=47\text{MeV}$	–	–	–	–

despite of enormous progress in the experiments no fully consistent picture as yet
 \Rightarrow further precision experiments needed: HADES: $p A \rightarrow \rho, \omega + X$ on bound proton

emerging evidence from all heavy-ion experiments (DLS, HADES, CERES, NA60)
 that low mass part of dilepton excess ($m_{l+l-} \leq 500 \text{ MeV}$) may be dominated by
 coupling to baryon resonances